

## Section A. Nomenclature Topics

### Part III. First Law of Thermodynamics

#### Chapter 9. Calculus Methods

### LECTURE 10. FUNCTIONS BY PFAFF

#### § 1.11. Application of calculus to thermodynamics

From the equations of

$$dq = du + dl. \quad (2.17)$$

$$dl = pdv. \quad (2.18)$$

[113, pp. 29-32]

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We get the first law of thermodynamics given in the view of

[113, p. 29, (2.25)]

$$dq = du + pdv. \quad (2.25)$$

From

$$dq = di - pdv - vdp + pdv = di - vdp, \quad (2.22)$$

[113, p. 29, (2.26)]

$$dq = di - vdp. \quad (2.26)$$

In mathematics they are **functions by Pfaff of two variable values**. The general view of these functions is [113, p. 29, (2.27)]

$$d\Phi = M(x, y)dx + N(x, y)dy. \quad (2.27)$$

for two and

$$d\Phi = X_1dx_1 + X_2dx_2 + X_3dx_3 + \dots + X_ndx_n.$$

for  $n$  – independent variables. Here  $x_1, x_2, x_3, \dots, x_n$  – independent variables, and  $M(x, y), N(x, y), X_1, X_2, X_3, \dots, X_n$  – functions of these variables.

For instance, pressure  $p$  and specific volume  $v$  are the independent variables, and then enthalpy  $i$  is a function of  $p$  and  $v$

$$i = f_{pv}(p, v).$$

That means that at some point with the parameters of  $p$  and  $v$  we will have a certain value of the function  $i$ . At the same time, for instance, work is not a function of the same independent variables, i.e. (that is) for some certain values of  $p$  and  $v$  one cannot tell what the work would be equaled to.

In mathematics there distinguished two kinds (types, sorts) of the functions by **Pfaff**. The **first kind** –  $d\Phi$  is the complete differential of some certain function of  $d\Phi(x, y)$ . In this case **the equality (condition by Euler) is realized** (implemented, fulfilled, carried into effect, accomplished, met) [113, p. 30, (2.28)]

$$\frac{\partial M(x, y)}{\partial y} = \frac{\partial N(x, y)}{\partial x}. \quad (2.28)$$

Such functions in mathematics are called **functions of point**, and in thermodynamics – **functions of state**.

For example, temperature is a function of a state of such independent parameters as pressure and specific volume. That means that in any state, which is characterized with the parameters of  $p$  and  $v$ , the temperature has quite a certain value, i.e. a function of  $T = f_T(p, v)$  exists.

The **second kind of the Pfaff functions** –  $d\Phi$  is not the complete differential of the function of  $d\Phi(x, y)$ . In this case **the equality (condition by Euler (2.28)) is not realized** (implemented, fulfilled, carried into effect, accomplished, met) [113, p. 30, (2.29)]

$$\frac{\partial M(x, y)}{\partial y} \neq \frac{\partial N(x, y)}{\partial x}. \quad (2.29)$$

The functions by Pfaff of the second kind are called functions of line in mathematics, and in thermodynamics – **functions of process**.

The complete differential of the function of  $\Phi$  is notated (written, put) down in the view of [113, p. 30, (2.30)]

$$d\Phi = \frac{\partial \Phi}{\partial x} dx + \frac{\partial \Phi}{\partial y} dy. \quad (2.30)$$

Let us show that if (2.30) is the complete differential, then the equation of (2.28) realizes.

Comparing (2.30) and (2.27), we find [113, p. 30, (2.31)]

$$M(x, y) = \frac{\partial \Phi}{\partial x}, \quad N(x, y) = \frac{\partial \Phi}{\partial y}. \quad (2.31)$$

After differentiating the first expression form (2.31) with respect to  $y$ , and the second with respect to  $x$ , we will get [113, p. 30, (2.32)]

$$\frac{\partial M(x, y)}{\partial y} = \frac{\partial^2 \Phi}{\partial x \partial y}, \quad \text{and} \quad \frac{\partial N(x, y)}{\partial x} = \frac{\partial^2 \Phi}{\partial y \partial x}. \quad (2.32)$$

From (2.32) we get (2.28). From here we can come to conclusion, that the condition of (2.28) is necessary for the existence of the complete differential of the function of  $\Phi(x, y)$ .

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In case if  $d\Phi$  is the complete differential of some function of  $F(x, y)$ , then the value of the integral does not depend upon the way of integration



For some independent parameters of  $x$  and  $y$  [113, p. 31, (2.33), (2.34)]

$$du = \frac{\partial u}{\partial x} dx + \frac{\partial u}{\partial y} dy, \quad (2.33)$$

$$dv = \frac{\partial v}{\partial x} dx + \frac{\partial v}{\partial y} dy. \quad (2.34)$$

Substituting (2.33) and (2.34) into (2.25), we will get

$$dq = \frac{\partial u}{\partial x} dx + \frac{\partial u}{\partial y} dy + p \left( \frac{\partial v}{\partial x} dx + \frac{\partial v}{\partial y} dy \right)$$

or

$$dq = M(x, y)dx + N(x, y)dy,$$

where [113, p. 32, (2.35), (2.36)]

$$M(x, y) = \frac{\partial u}{\partial x} + p \frac{\partial v}{\partial x}, \quad (2.35)$$

$$N(x, y) = \frac{\partial u}{\partial y} + p \frac{\partial v}{\partial y}. \quad (2.36)$$

Differentiating (2.35) and (2.36) with respect to  $y$  and  $x$  correspondingly

$$\frac{\partial M(x, y)}{\partial y} = \frac{\partial^2 u}{\partial x \partial y} + \frac{\partial p}{\partial y} \frac{\partial v}{\partial x} + p \frac{\partial^2 v}{\partial x \partial y}, \quad (2.35a)$$

$$\frac{N(x, y)}{\partial x} = \frac{\partial^2 u}{\partial y \partial x} + \frac{\partial p}{\partial x} \frac{\partial v}{\partial y} + p \frac{\partial^2 v}{\partial y \partial x}. \quad (2.36a)$$

Making transformations

$$\frac{\partial M(x, y)}{\partial y} - \frac{N(x, y)}{\partial x} = \frac{\partial^2 u}{\partial x \partial y} + \frac{\partial p}{\partial y} \frac{\partial v}{\partial x} + p \frac{\partial^2 v}{\partial x \partial y} - \left( \frac{\partial^2 u}{\partial y \partial x} + \frac{\partial p}{\partial x} \frac{\partial v}{\partial y} + p \frac{\partial^2 v}{\partial y \partial x} \right).$$

After rearranging

$$\frac{\partial M(x, y)}{\partial y} - \frac{N(x, y)}{\partial x} = \frac{\partial p}{\partial y} \frac{\partial v}{\partial x} - \frac{\partial p}{\partial x} \frac{\partial v}{\partial y} = D.$$

This is the **determinant by Jacoby (Jacobean)**

$$D = \begin{vmatrix} \frac{\partial v}{\partial x} & \frac{\partial p}{\partial x} \\ \frac{\partial v}{\partial y} & \frac{\partial p}{\partial y} \end{vmatrix}.$$

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[113, p. 29-32]